



# Quasiparticle spectrum and optical properties of SnO<sub>2</sub>

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# Outline

- Introduction
- Electronic structure
- Optical absorption
- Quasiparticle spectrum of  $\text{SnO}_2$
- Pressure effects
- Scalar relativistic pseudopotential
- Optical properties
- Conclusions and outlook



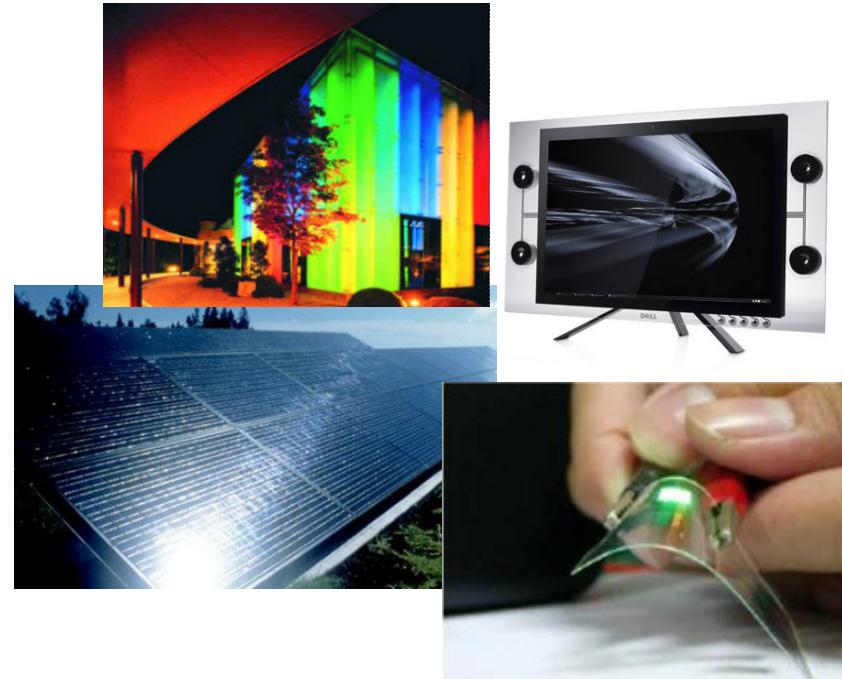
# Introduction

Transparent  
conducting oxides:

- band gap  $\geq 2$  eV
- resistivity  $\leq 10^{-4} \Omega \text{ cm}$   
and  $\mu \geq 50 \text{ cm}^2/\text{V s}$

$\text{SnO}_2$ :

- band gap  $\sim 3.6$  eV<sup>1</sup>
- Ta-doped<sup>2</sup> resistivity  
 $\sim 2 \times 10^{-4} \Omega \text{ cm} \dots$   
and  $\mu \sim 60 \text{ cm}^2/\text{V s}$



<sup>1</sup>M. Nagasawa and S. Shionoya, Phys. Lett. 22, 409 (1966).

<sup>2</sup>S. Nakao *et al.*, Appl. Phys. Express 3, 031102 (2010)



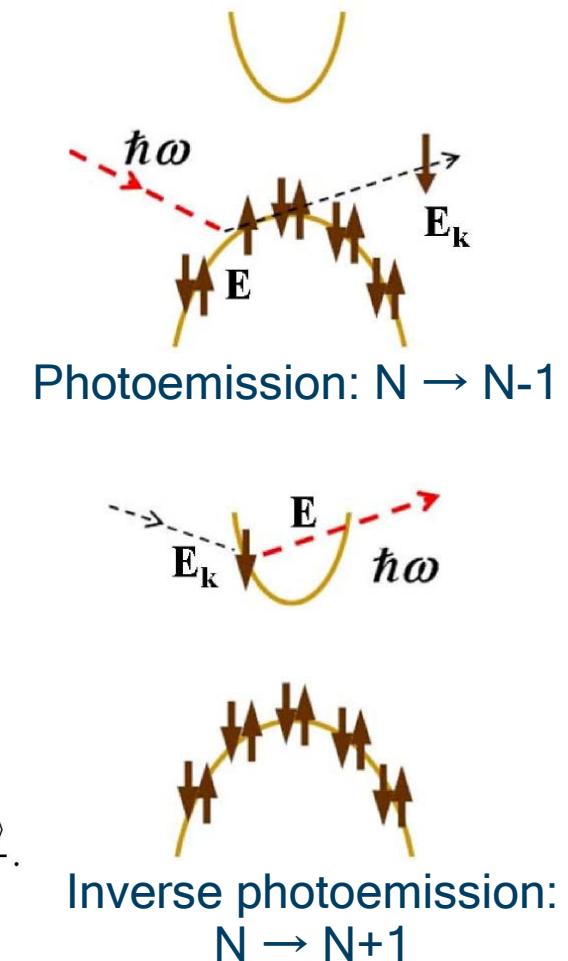
# Electronic structure

The electron Green function describes the evolution of an excitation

$$G(\mathbf{r}_1, t_1; \mathbf{r}_2, t_2) = -\frac{i}{\hbar} \langle \Psi_0 | T[\hat{\psi}(\mathbf{r}_1, t_1) \hat{\psi}^\dagger(\mathbf{r}_2, t_2)] | \Psi_0 \rangle$$

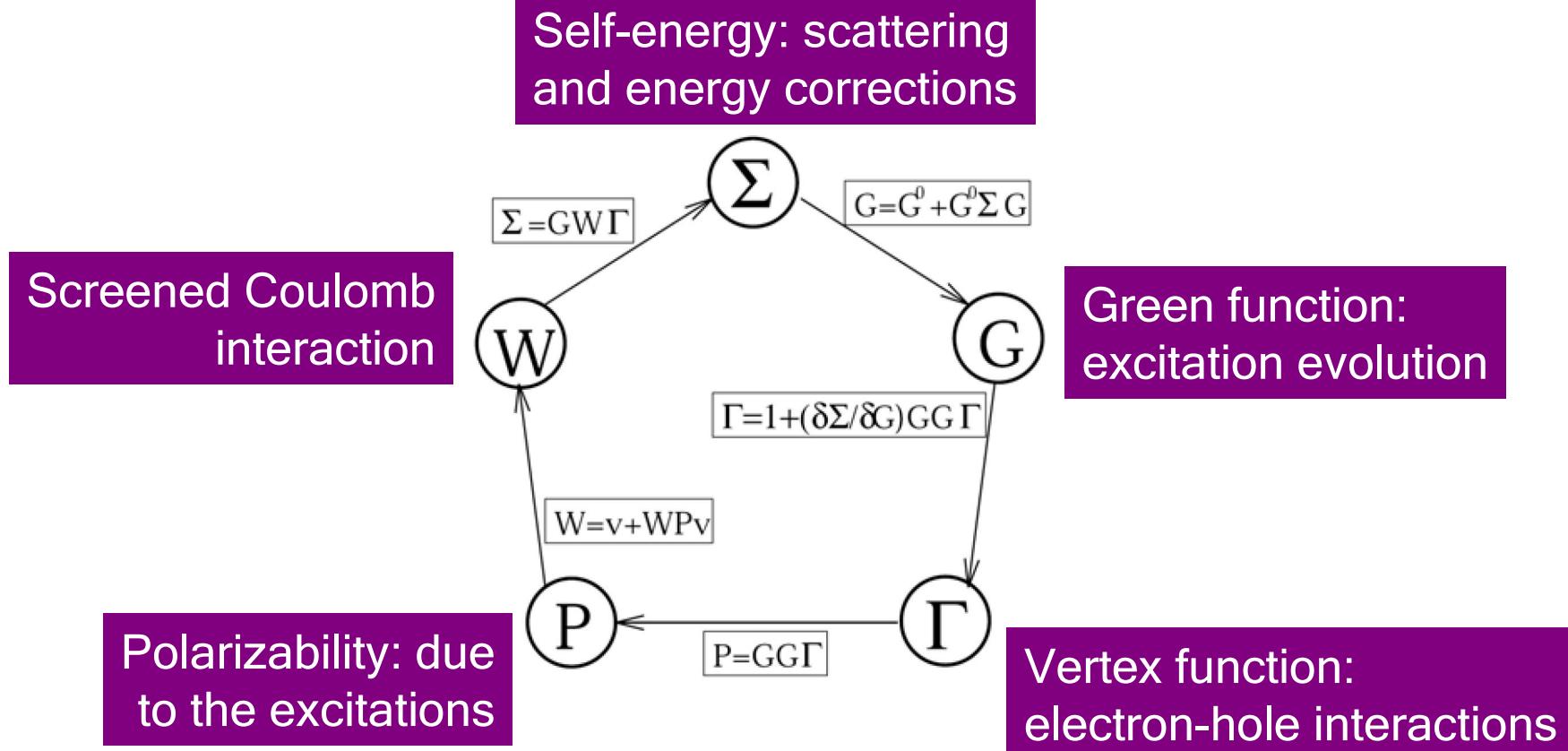
The Green function poles are single-particle excitation energies

$$\begin{aligned} G(\mathbf{r}_1, \mathbf{r}_2; \omega) &= \frac{1}{\hbar} \sum_i \frac{\langle \Psi_0^{(N)} | \hat{\psi}(\mathbf{r}) | \Psi_i^{(N+1)} \rangle \langle \Psi_i^{(N+1)} | \hat{\psi}^\dagger(\mathbf{r}') | \Psi_0^{(N)} \rangle}{\hbar\omega - \epsilon_i^{(N+1)} + i\eta} \\ &+ \frac{1}{\hbar} \sum_i \frac{\langle \Psi_0^{(N)} | \hat{\psi}^\dagger(\mathbf{r}') | \Psi_i^{(N-1)} \rangle \langle \Psi_i^{(N-1)} | \hat{\psi}(\mathbf{r}) | \Psi_0^{(N)} \rangle}{\hbar\omega - \epsilon_i^{(N-1)} - i\eta}. \end{aligned}$$





## Hedin equations & GW approximation<sup>1,2</sup>



<sup>1</sup>L. Hedin, Phys. Rev. **139**, A796 (1965).

<sup>2</sup>W. G. Aulbur, L. Jönsson, and J. W. Wilkins, Solid State Phys. **54**, 1 (1999).



## GW scheme

From the Dyson equation

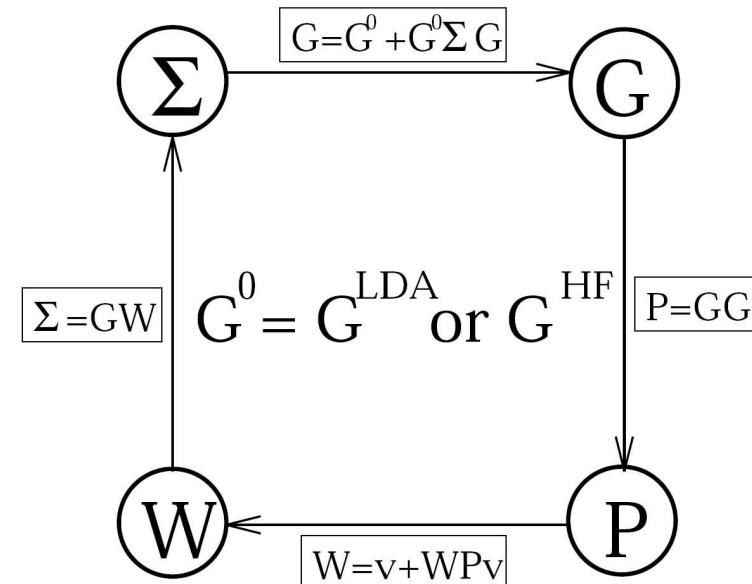
$$\mathcal{H}_0(\mathbf{r})\psi_j(\mathbf{r}) + \int d^3r' \hbar \Sigma^*(\mathbf{r}, \mathbf{r}'; \epsilon_j) \psi_j(\mathbf{r}') = \epsilon_j \psi_j(\mathbf{r}).$$

DFT-local density approx.

$$\Sigma(\mathbf{r}, \mathbf{r}'; \omega) = \delta(\mathbf{r} - \mathbf{r}') V_{\text{xc}}(\mathbf{r})$$

“GW” approximation

$$\Gamma = 1 \Rightarrow \Sigma = \text{GW}$$

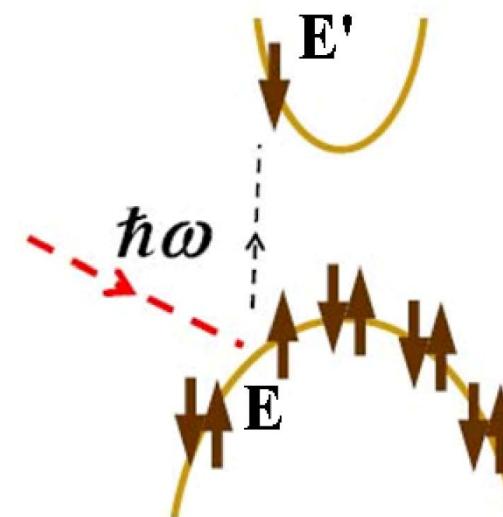


# Optical absorption

- Macroscopic dielectric function

$$\begin{aligned}\varepsilon_M(\omega) &\equiv \lim_{\mathbf{q} \rightarrow 0} \frac{1}{\varepsilon_{\mathbf{G}=0, \mathbf{G}'=0}(\mathbf{q}, \omega)}, \\ &= 1 - \lim_{\mathbf{q} \rightarrow 0} [v(\mathbf{q})_0 \bar{P}_{\mathbf{G}=\mathbf{G}'=0}(\mathbf{q}, \omega)]\end{aligned}$$

Photoabsorption: two-particle excitation process ...  
but the GW polarizability does not contain electron-hole interactions!

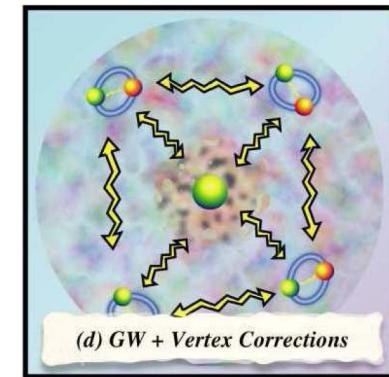




- Bethe-Salpeter

Vertex corrections:

$$\frac{\delta \Sigma(12)}{\delta G(45)} \simeq i\hbar \delta(14)\delta(25)W(1^{+2})$$



1

=> Bethe-Salpeter equation for the polarisability

$${}^4\overline{P} = {}^4P_{IQP} + {}^4P_{IQP}K\ {}^4\overline{P},$$

$$K(1234) = \delta(12)\delta(34)\overline{v}(13) - \delta(13)\delta(24)W(12)$$

${}^4P$  = “4-point” polarisability

${}^4P_{IQP}$  = independent quasiparticle polarisability

$v$  = bare Coulomb interaction

$W$  = dynamically screened Coulomb interaction



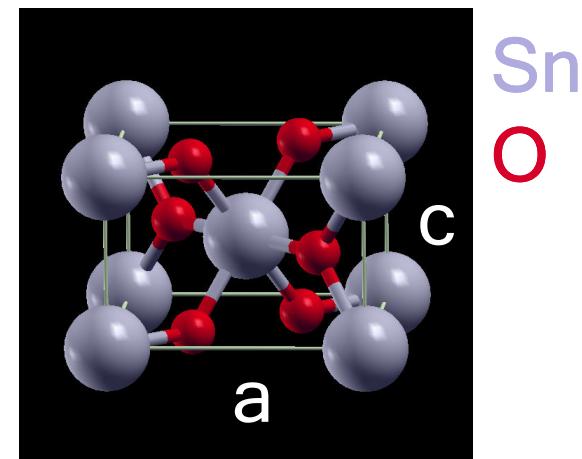
# Quasiparticle spectrum of SnO<sub>2</sub>

- DFT calculations

22e PP required: [Ge]4s<sup>2</sup>4p<sup>6</sup>4d<sup>10</sup>5p<sup>2</sup>5s<sup>2</sup>  
(Opium code<sup>1</sup>) ... 105 Ha ecut!

## Structural properties

	a (Å)	c (Å)	u	B (GPa)
Exp. <sup>2</sup>	4.7374	3.1864	0.30562	212.3 <sup>(3)</sup>
Theory	4.7154	3.1864	0.30605	211.7



SnO<sub>2</sub> structure:  
Rutile (*P4<sub>2</sub>mnm*)

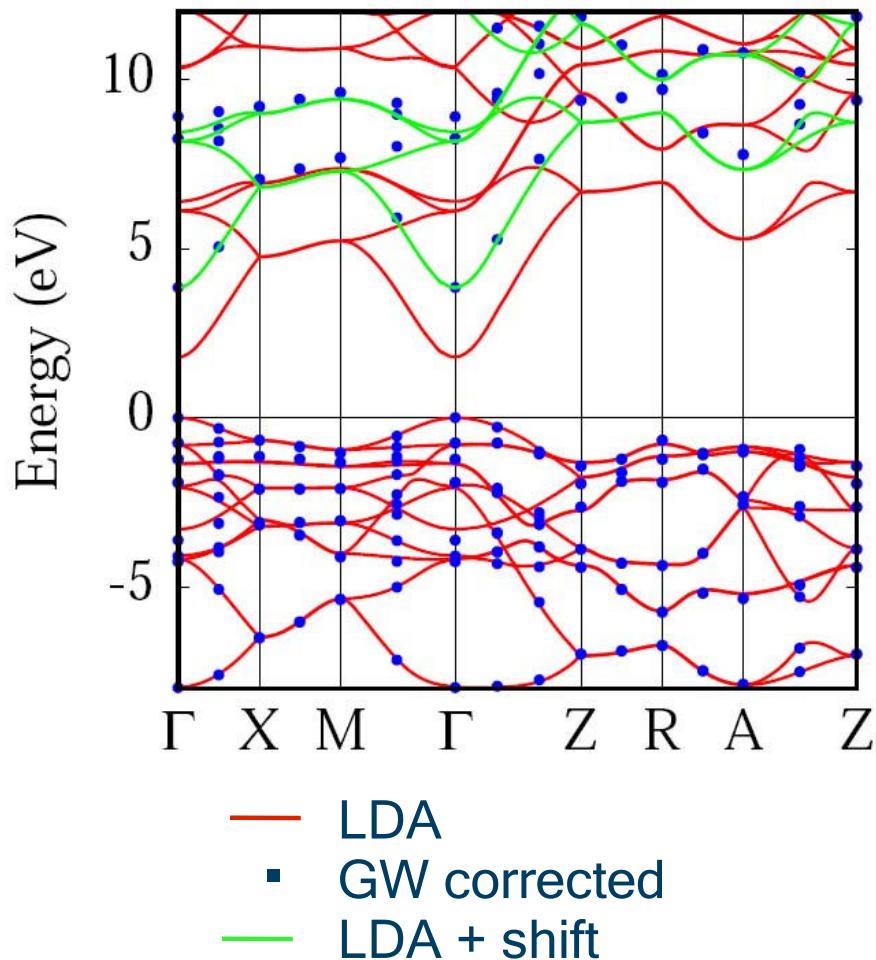
<sup>1</sup>[opium.sourceforge.net/index.html](http://opium.sourceforge.net/index.html)

<sup>2</sup>A. Bolzan *et al.*, Acta Cryst. B53, 373 (1997).

<sup>3</sup>E. Chang and E. Graham, J. Geophys. Res. 80, 2595 (1975).



## • Electronic structure



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For ref., optical gap at  $\Gamma$ : 3.6 eV

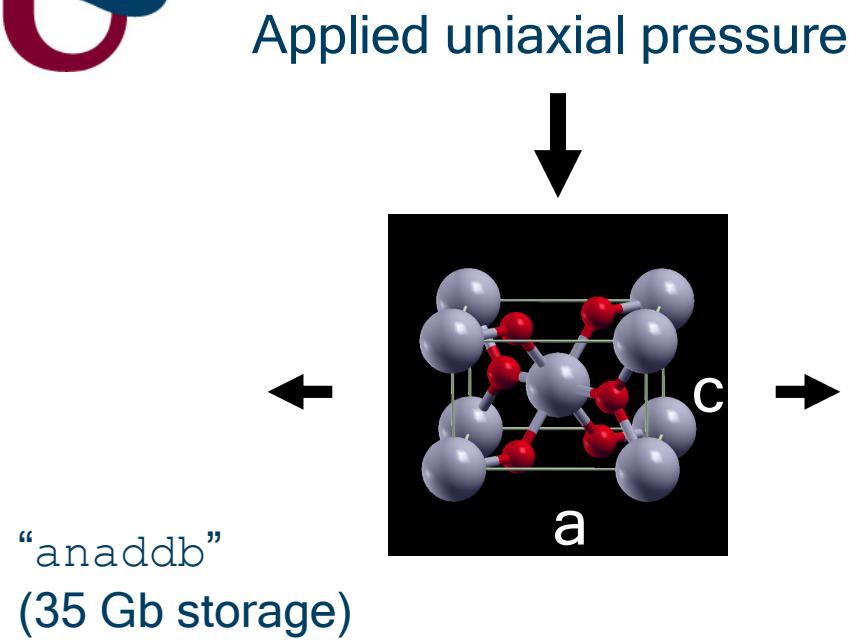
Gap values (eV):

k pt	$\Gamma$	X	M	Z
LDA	1.80	5.42	6.17	7.99
GW	3.85	7.69	8.72	10.81
$\Delta E$	2.05	2.27	2.55	2.82

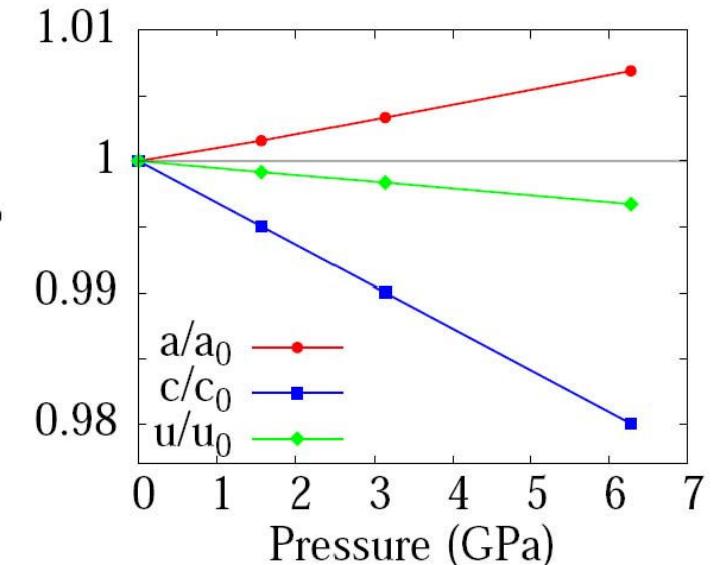
Effective masses ( $m_e$ ):

	$m^*_\perp$	$m^*_\parallel$	$m^*_p$
Exp <sup>1</sup> .	0.299	0.234	0.275
Theory	0.253	0.223	0.271

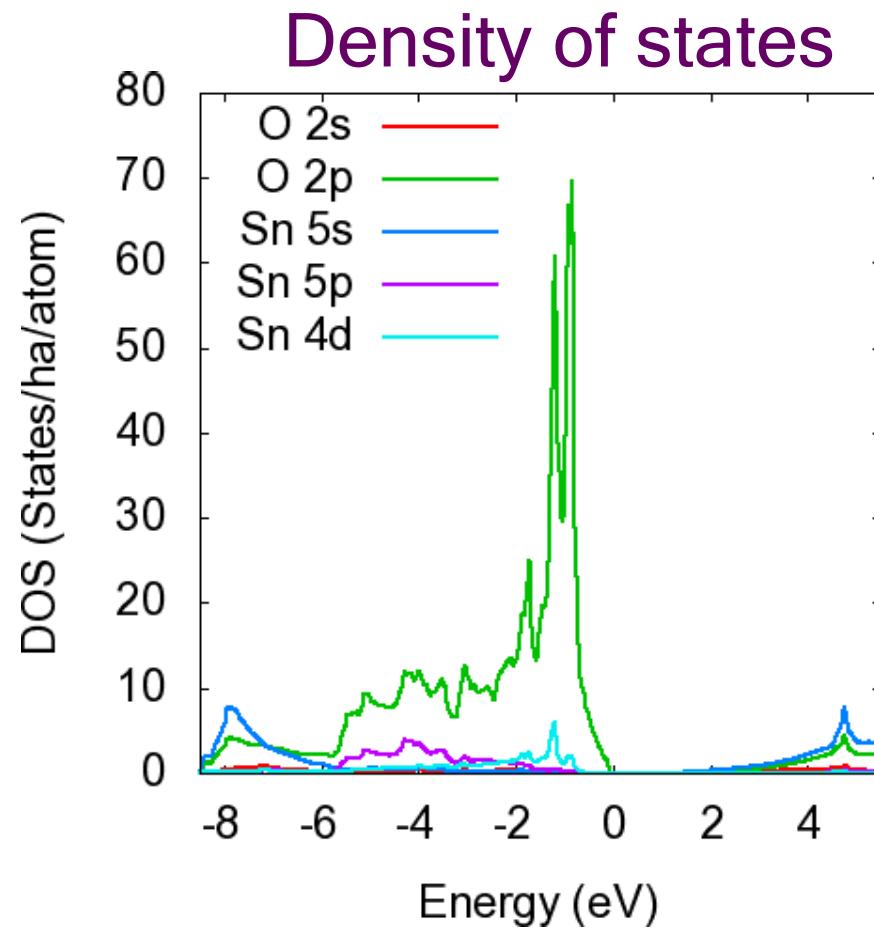
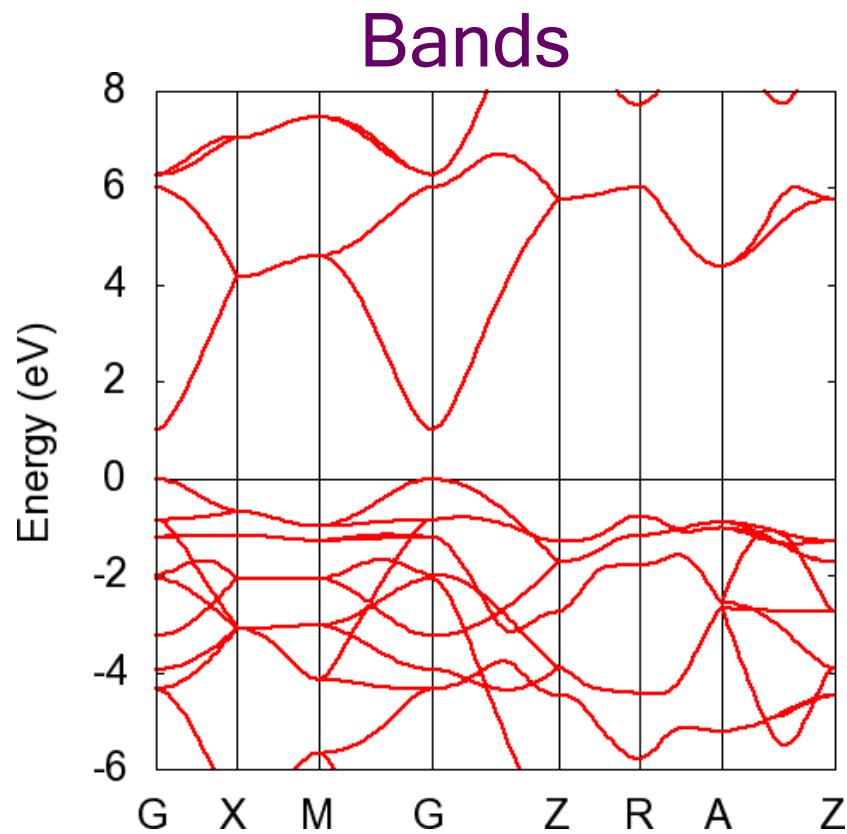
<sup>1</sup>K. Button et al., Phys. Rev. B 4, 4539 (1971).



## Pressure effects



	$C_{11}$	$C_{33}$	$C_{12}$	$C_{13}$	$C_{44}$	$C_{66}$	GPa
Exp. <sup>1</sup>	261.7	449.6	177.2	155.5	103.1	207.4	
Theory	274.3	412.5	180.9	149.8	94.3	202.9	
	$S_{11}$	$S_{33}$	$S_{12}$	$S_{13}$	$S_{44}$	$S_{66}$	$\text{TPa}^{-1}$
Theory	6.801	3.186	-3.913	-1.049	10.608	4.929	

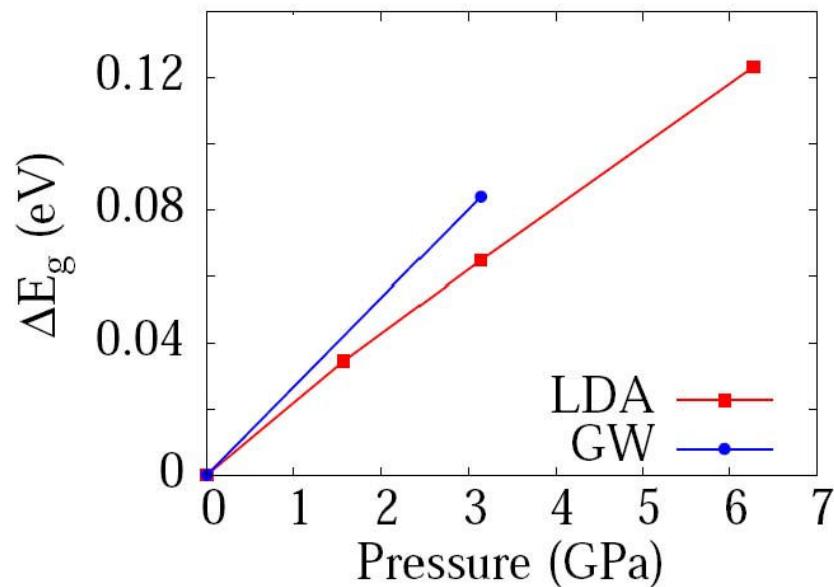




## Pressure and GW

GW corrections at high symmetry k-points

k pt	$\Gamma$	X	M	Z
$\Delta E (P=0)$	2.05	2.27	2.55	2.82
$\Delta E (P=3.1 \text{ GPa})$	2.08	2.30	2.57	2.84



$$P_{\text{coeff}} \text{ GW} = 27 \text{ meV/GPa}$$

$$P_{\text{coeff}} \text{ GW}/P_{\text{coeff}} \text{ LDA} = 0.74$$



# Scalar relativistic pseudopotential

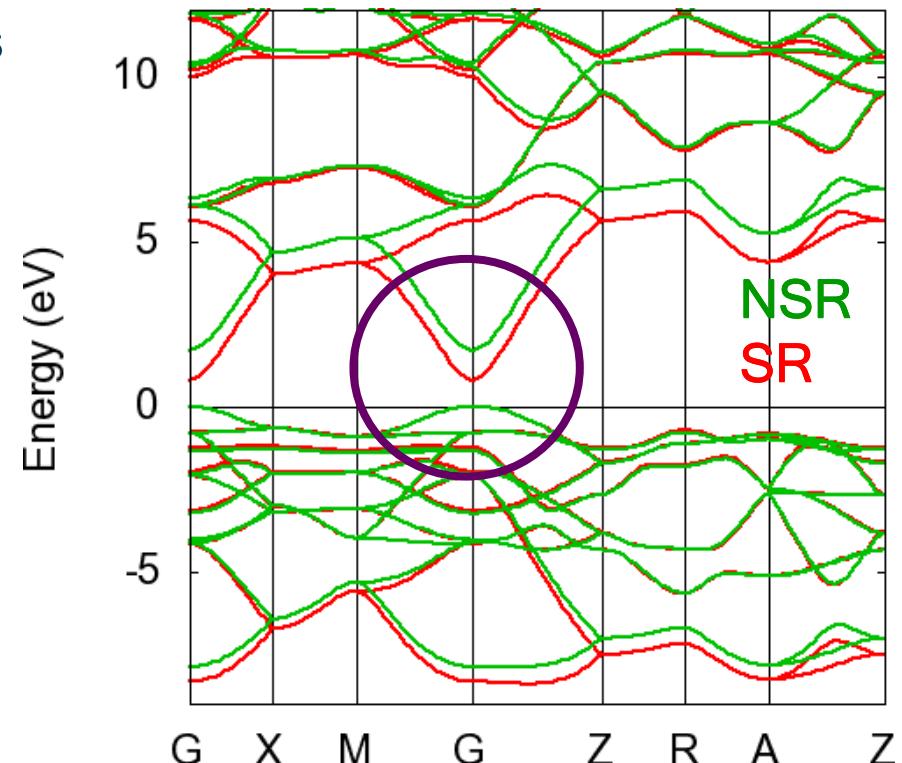
Scalar relativistic effects

- ⇒ shrinkage of core s and p shells
- ⇒ change of bandwidths and  $E_g$

Structural parameters

	a (Å)	c (Å)	u
Exp. <sup>2</sup>	4.7374	3.1864	0.30562
Theory	4.7185	3.1849	0.30620

SR  $E_g = 0.81$  eV  
(vs 1.80 eV in NSR !)





## Elastic constants

	$c_{11}$	$c_{33}$	$c_{12}$	$c_{13}$	$c_{44}$	$c_{66}$
Exp. <sup>1</sup>	261.7	449.6	177.2	155.5	103.1	207.4
Theory	238.7	416.9	177.4	152.8	91.6	204.4
	$s_{11}$	$s_{33}$	$s_{12}$	$s_{13}$	$s_{44}$	$s_{66}$
Theory	9.808	3.281	-6.520	-1.205	10.916	4.892

$$B_{SR} = 202.6 \text{ GPa}$$
$$B_{NSR} = 211.7 \text{ GPa}$$
$$B_{exp} = 212.3 \text{ GPa}$$

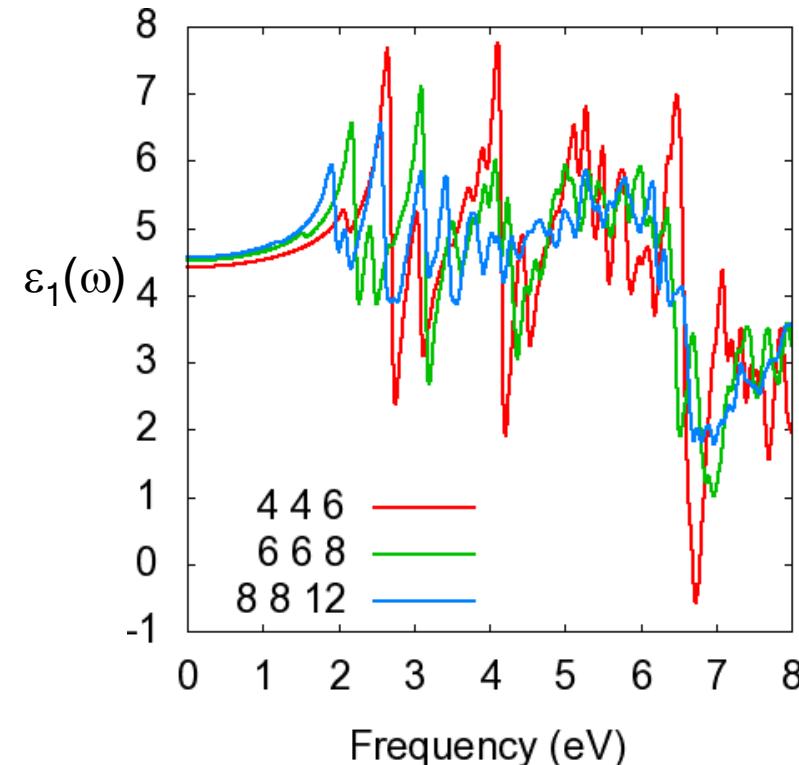
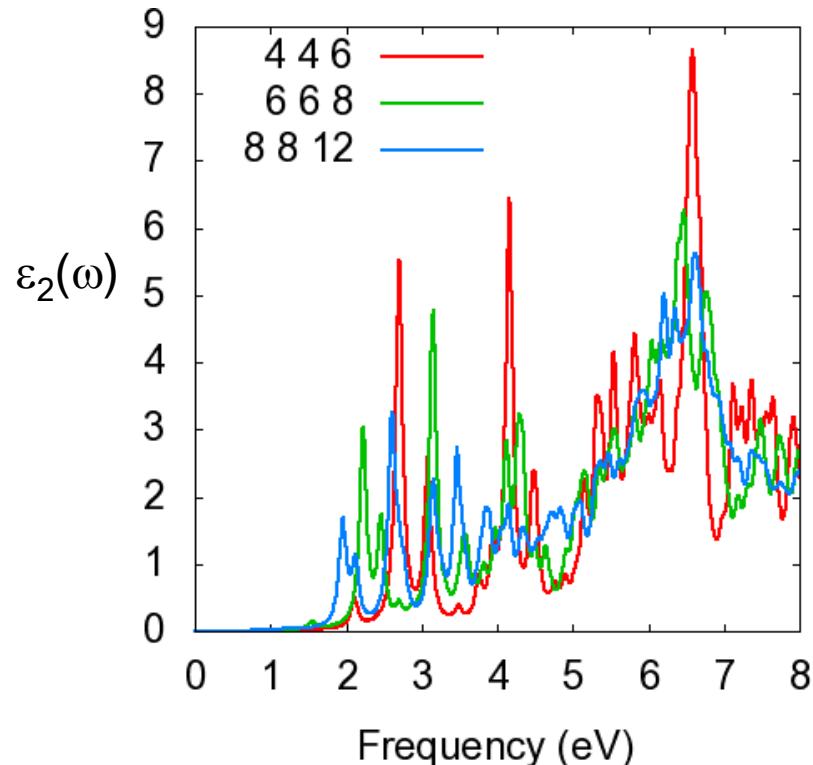
GW gaps (eV)					Effective masses ( $m_e$ )			
k pt	$\Gamma$	X	M	Z		$m^*_\perp$	$m^*_\parallel$	$m^*_p$
NSR	3.85	7.69	8.72	10.81	NSR	0.253	0.223	0.271
SR	2.65	6.75	7.59	9.53	SR	0.189	0.170	0.204

Gaps go down

... and effective masses go down



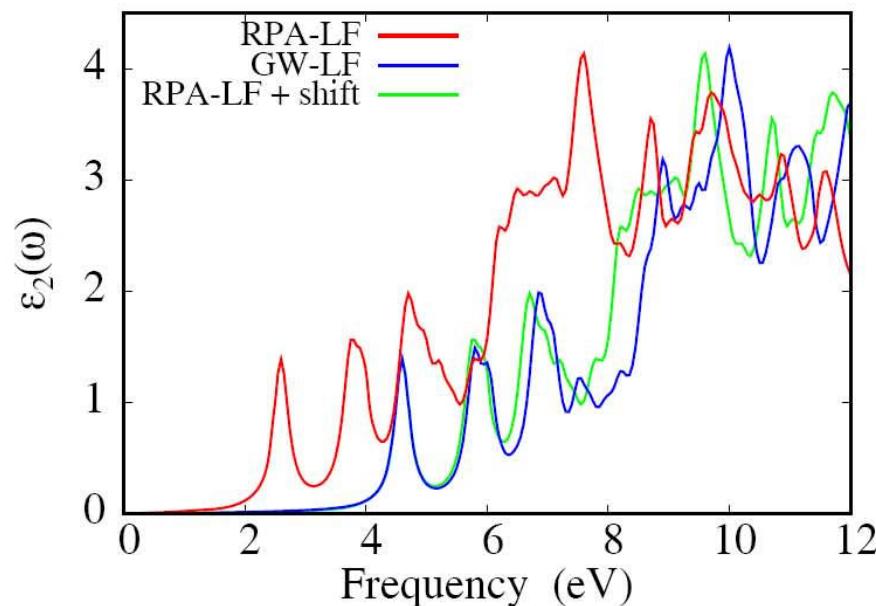
“optic” code



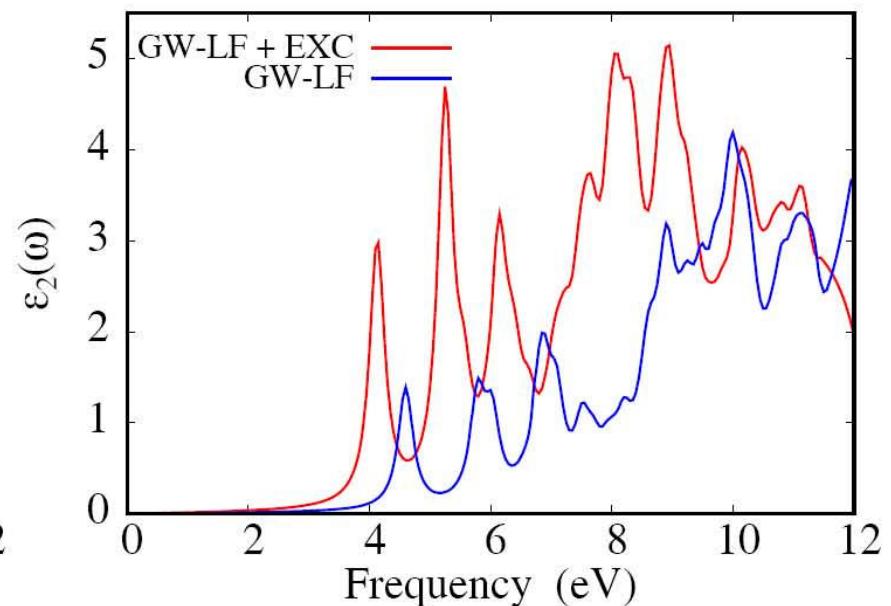
Caution: dense k-point needed ( $> 5\text{Gb}$  memory/proc.)



## Imaginary part of the macroscopic dielectric function: LDA-RPA vs GW-RPA vs Bethe-Salpeter



Absorption edge and structure  
at higher energies



Absorption edge and structure  
also at low energies



## Conclusions and outlook

- Structural and elastic properties are accurately predicted.
- GW corrections are band and momentum dependent.
- ... but require more work: quasi-particle self-consistent calculation (cassiterite\_i\_SCR = 65 Gb, GW corrections = 11 Gb/proc. with spectral method).
- Bethe-Salpeter shows important excitonic effects.  
Future work: Wannier interpolation for GW eigenvalues and LDA for wavefunctions?